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# **Increasing Electrical Damping in Energy Harnessing Transducers**

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#### Abstract

Energy and power in tiny batteries are often insufficient to sustain the demands of a wireless microsystem for extended periods. The transducers are viable alternatives because they draw power from a vast tank-free supply of ambient kinetic energy in vibrations. Unfortunately, small devices alone seldom dampen vibrations enough to fully harness what is available, which is why investing energy to increase the electrical damping force that transducers impose is so important. This paper introduces and evaluates three investment schemes and 0.35-µm CMOS switchedinductor circuits that increase this force to generate more output power.

#### I. Introduction

Energy from ambient sources (as in Fig. 1) can extend the operational life of a microsystem by recharging a battery. State-of-the-art depleting microscale transducers, however, only generate µW's, of which power conditioning circuits consume a portion [2]. Fortunately, electrical energy  $E_E$  increases with electrical damping force, which as this paper demonstrates, can increase with initially invested energy E<sub>INV</sub>. To consider this in more detail, Sections II, III, IV and V discuss how investing energy electromagnetic, increases output power in electrostatic, and piezoelectric transducers, drawing relevant conclusions in Section VI.



Fig. 1. Sample harvesting wireless microsystem.

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### II. Electrical Damping A. Theoretical Behavior

One such alternative for increased damping is to introduce damping on the electrical side rather than the mechanical side of the system's drive motor. Using a one degree-of-freedom device modeling a virtual wall as an example, it is seen that this can be accomplished by inserting an electrical resistance in parallel with the motor (Fig. 1b) [3-5]. After analyzing the constitutive relations of such a system it is seen that the equivalent mechanical damping that this electrical system adds is

$$B_{eq} = \frac{K_t^2}{\left(R_1 + R_m\right)}.$$
(1)

where Kt is the motor's torque constant, Rm is the motor's internal resistance, and R1 is the added parallel resistance.



Figure 2: A mechanically damped system (a) and two electrically damped systems; one without (b) and one with frequency dependence (c).

With only an added resistor the electrical system acts just as a mechanical damper, dissipating energy throughout the device's range of motion.



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An improvement can be made, however, by adding a capacitor in series with the added parallel resistance (Fig. 1c). This makes the electrical damping frequency dependent. The values of resistance and capacitance can be chosen to give the system a cutoff frequency around the normal bandwidth of human hand motion (a relatively small 2 Hz [6-8] to 4.5 Hz [9]). Thus, away from any constraints, when movement is governed almost entirely by inputs from the human user, the system acts as if there is no extra damping. When a high frequency event occurs, as in impacting a virtual wall, the electrical damping can serve to prevent the energy growth that leads to limit cycle oscillations and other instabilities. This method of providing real physical damping, therefore, eliminates the need for unwieldy mechanical dampers while also simplifying the control structure and device design by doing away with the need for negative virtual damTop ifnugr.t her understand the behavior of an electrically damped system, a one degree-of-freedom device with electrical damping can be modeled (as in Fig. 2) and the following system transfer function can be analyzed. This model omits friction and structural compliance for simplicity.

 $\frac{1}{LCs^2 + (R_1 + R_m)Cs + 1}$ 

 $\tau(s) = A(s)I(s) - Z(s)v(s)$ 

Motor

I. K<sub>t</sub>

 $R_1 = C$ 

....I

Figure 3: Model used for analyzing the theoretical

behavior of a one degree-of-freedom haptic display

with electrical damping

ന

 $\frac{K_t^2 n^2 Cs}{+(R_t + R_m)Cs + 1} + n^2 (B + Js) v(s)$ 

ser in nut2

(2a)

(2b)

Here, it is seen that the torque,  $\tau$ , is responsive to two inputs: the current from the amplifier, I, and the angular velocity of the motor shaft, v. The system characteristics of the device used in testing (described in Section 3) can be substituted into (2a) and the resulting frequency responses can be plotted to obtain more complete picture of haptic a display performance. First, velocity is assumed to be zero and the resulting plot of the magnitude of A(s), Fig. 3a, shows the frequency response of torque to a current input (specific parameter values are given in Section 3). It is desirable to have this plot constant, or as close as possible, because any shift in this effective "torque constant" corresponds [10] to a change in the ability of a commanded current to output a desired torque. While the goal of electrical damping is to dissipate unwanted energy at high frequencies, the ability to control the haptic display with current commands of reasonable magnitude, at all frequencies, must be maintained. Magnitude of Frequency Response for A(s), the Torque/Current Transfer Function for Electrical Damping = 0.00755Nms/rad



Figure 4: Magnitude portions of the Bode plots for the transfer functions A(s) (a) and Z(s) (b) for a system with 0.00755 Nm/(s/rad) of electrical damping

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If current rather than velocity is assumed zero, the magnitude of Z(s), the transfer function from velocity input to torque output can be plotted (Fig. 3b). This magnitude can be further broken down into its real and imaginary components (see Fig. 4).  $Re{Z(s)}$ corresponds to the effective damping of the system while  $Im{Z(s)}/\omega$  represents the apparent inertia felt at the device output [[11-13]. From these plots it is clear that significant additional damping is added to the system at high frequencies, and only high frequencies. Also, motions at low frequencies experience only a slight increase in system inertia due to the added parallel capacitance. Thus, electrical damping can aid in stabilizing high frequency events like an impact with a virtual wall, while not hindering a user's unconstrained motion away from the boundary. Effective Damping vs. Frequencfor Electrical Damping = 0.00755 Nms/rad



(a) Apparent Inertia vs. Frequency (b) Figure 5: Theoretical effective damping (a) and apparent inertia (b) obtained from Z(s) for a system with 0.00755 Nm/(s/rad) of electrical damping

#### II. ELECTROSTATIC TRANSDUCERS

A variable capacitor  $C_{VAR}$  having one physically suspended plate that moves under the influence of environmental motion can harvest energy. In voltageconstrained (VC) harvesting, because capacitancevoltage product represents charge, maintaining the voltage across  $C_{VAR}$  constant when vibrations separate its plates (i.e., decrease capacitance) reduces its charge, which means  $C_{VAR}$  produces energy. In charge constrained (QC) operation, since linear variations in  $C_{VAR}$ 's voltage causes squared changes in energy (i.e.,  $E_C$  is  $0.5C_{VAR}v_C^2$ ), fixing  $C_{VAR}$ 's charge by keeping it open when vibrations decrease  $C_{VAR}$  raises  $v_C$ , so  $v_C^2$ increases surpass linear reductions in  $C_{VAR}$  to produce a net energy gain.

#### **A. Increasing Electrical Damping**

In both VC and QC operation, the system invests energy at the beginning of each cycle to pre-charge C<sub>VAR</sub>. Some or all of this charge remains on C<sub>VAR</sub>'s plates through the harvesting phase to establish an electrostatic attraction that opposes (and damps) the physical movement of the suspended plate. Vibrations, as a result, produce more energy when this electrical damping force  $(F_{DE})$  is higher [14-15]. In the presence of overpowering mechanical damping forces (when Z<sub>s</sub> overwhelms  $k_C^2 Z_E$  in Fig. 2),  $F_{DE}$  has little impact on the displacement x(t) of  $C_{VAR}$ 's plates [4], which means raising F<sub>DE</sub> draws more electrical energy from vibrations. Therefore, because F<sub>DE</sub> increases with the square of C<sub>VAR</sub>'s voltage v<sub>C</sub>, as does C<sub>VAR</sub>'s E<sub>C</sub>, higher voltages through the harvesting phase induce more electrical damping in the transducer and, as a result, produce more output energy E<sub>H</sub>:

$$\mathbf{E}_{\mathrm{H}} = \int \mathbf{F}_{\mathrm{DE}} \mathrm{d}\mathbf{x} \propto \int \frac{\mathbf{v}_{\mathrm{C}}^{2}}{\mathbf{x}(\mathrm{t})^{2}} \mathrm{d}\mathbf{t} \propto \int \frac{\mathbf{E}_{\mathrm{C}}}{\mathbf{x}(\mathrm{t})^{2}} \mathrm{d}\mathbf{t} \,.$$

This means that keeping  $v_C$  as close to  $C_{VAR}$ 's breakdown voltage ( $V_{MAX}$ ) throughout the harvesting period generates more energy than otherwise, which is why VC harvesting at or near  $V_{MAX}$  spawns more energy than in QC operation, where  $v_C$  rises and nears  $V_{MAX}$  only at the end of the cycle [9].



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#### **B.** Voltage-clamping Capacitor

Constraining  $C_{VAR}$ 's voltage with a 2.7 – 4.2-V li-ion battery [10] through the harvesting cycle is one way of extracting energy from motion directly into a battery (V<sub>BAT</sub>). The advantages of this are that no additional capacitors or energy transfers, which are lossy, are necessary. Unfortunately, V<sub>BAT</sub> is not the maximum voltage C<sub>VAR</sub> can sustain, which means C<sub>VAR</sub> does not draw as much energy as its breakdown voltage allows. So, at the cost of silicon or printed-circuit-board (PCB) area, a large clamping capacitor C<sub>CLAMP</sub> (of up to 1 nF) that constrains C<sub>VAR</sub> (e.g., 50 – 250 pF) above V<sub>BAT</sub> near V<sub>MAX</sub>, as in Fig. 6, can harness sufficient energy to overcome losses in an additional energy-transfer phase.



Fig. 6. Constraining C<sub>VAR</sub>'s voltage with a clamping capacitor.

#### **Permanent Connection:**

In hard-wiring  $C_{CLAMP}$  to  $C_{VAR}$  [9], energy-transfer circuit (i.e.,  $T_X$ ) first invests energy  $E_{INV}$  from battery  $V_{BAT}$  to pre-charge both  $C_{VAR}$  and  $C_{CLAMP}$  close to  $V_{MAX}$ . Then, once the harvesting cycle ends,  $T_X$  must fully discharge both capacitors, before  $C_{VAR}$  uses remnant energy to help pull its plates together. Because  $C_{CLAMP}$  is much higher than  $C_{VAR}$  (to ensure  $C_{CLAMP}$  clamps  $C_{VAR}$  near  $V_{MAX}$  when  $C_{VAR}$  changes),  $T_X$  transfers considerably more energy ( $E_{INV}$  and  $E_H$ ) than it harvests ( $E_H$ ), so conduction losses are correspondingly higher.





#### **Asynchronous Connection:**

 $T_X$  in [11] pre-charges  $C_{VAR}$  to a fraction of  $V_{MAX}$  so mechanical energy can raise  $C_{VAR}$ 's voltage to a diode voltage above  $C_{CLAMP}$ 's initially high voltage (near  $V_{MAX}$ ) before driving charge into  $C_{CLAMP}$ .

Volume No: 4 (2017), Issue No: 3 (March) www.ijmetmr.com The interface circuit then transfers harnessed energy in  $C_{CLAMP}$  to  $V_{BAT}$ . Although  $T_X$  transfers less energy because  $C_{CLAMP}$  keeps its initial charge, the diode dissipates power and  $C_{VAR}$ 's voltage is considerably below  $V_{MAX}$  for a substantial portion of the harvesting period.

#### Managed Connection [Proposed]:

Alternatively,  $T_X$  in Fig. 7 charges  $C_{VAR}$  to  $C_{CLAMP}$ 's initial voltage (near  $V_{MAX}$ ), and once done, the controller closes switch  $S_3$  to steer mechanical energy extracted into  $C_{CLAMP}$ .  $T_X$  then discharges  $C_{VAR}$  into  $V_{BAT}$  before vibrations push its plates together, and deenergizes  $C_{CLAMP}$  with  $C_{VAR}$  (via  $S_3$ ) less often, when  $C_{CLAMP}$  reaches  $V_{MAX}$ . As such,  $C_{VAR}$  remains close to  $V_{MAX}$  through the entire harvesting phase and  $S_3$ dissipates less power than the diode in [11] (because its terminal voltages are considerably lower). Adding intelligence to manage the precharge process and the ensuing connection this way, however, requires energy, which represents a loss to the system.

#### **C. Performance and Limitations**

The major drawback to C<sub>CLAMP</sub> is its impact on integration. Unfortunately, reducing capacitance increases C<sub>CLAMP</sub>'s voltage variation (through the harvesting phase), so its voltage must start further below  $V_{MAX}$  (at  $V_{INI}$  in Fig. 8) to keep  $C_{CLAMP}$  from breaking down. As a result, C<sub>VAR</sub> harvests less energy per cycle, as E<sub>OUT</sub> in Fig. 8 shows below 100 pF for a 0.35-µm CMOS circuit with 40-V devices. Interestingly, increasing C<sub>CLAMP</sub> when permanently connected to C<sub>VAR</sub> (e.g., above 100 pF in Fig. 8) does not always increase  $E_{OUT}$ . This happens because  $T_X$ transfers more charge to raise C<sub>CLAMP</sub> near V<sub>MAX</sub>, which means additional conduction losses negate the gains of increased electrical damping forces. The circuit proposed in Fig. 7, however, transfers substantially less energy because C<sub>CLAMP</sub> retains its initial charge through all phases.



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Fig. 8. Simulated output energy across clamping capacitances.

One difference between the two connection strategies is the presence of  $S_3$ , which requires silicon area. Removing  $S_3$  and dedicating its area to other switches decreases the resistance across (and conduction losses in) the system, raising  $E_{OUT}$ . Reducing resistances by this amount, however, does not compensate for the losses that transferring all of  $C_{CLAMP}$ 's charge incurs, as  $E_{OUT}$  in Fig. 8 shows. Still, controlling  $S_3$  requires quiescent and switching energy not accounted for in Fig. 8. As a result, managing the connection is better only if conduction losses with a permanent connection exceed controller losses, which is more likely when  $C_{CLAMP}$  is higher because higher capacitance requires more energy to charge.

#### IV. PIEZOELECTRIC TRANSDUCERS A. Battery-coupled Damping

Piezoelectric transducers (PZT) generate charge in response to mechanical vibrations. When opencircuited, the resulting current energizes and deenergizes the capacitance across the surfaces of the device ( $C_P$ ) and supplies the parasitic leakage across the same (via  $R_P$ ). Cascading a full-wave rectifier and a battery  $V_{BAT}$  (as in Fig. 9, but without  $S_{RE}$  and  $L_{RE}$ ) steers charge away from  $C_P$  into  $V_{BAT}$  when PZT current  $i_P$  charges  $C_P$  above the barrier voltage that conducting diodes and  $V_{BAT}$  establish (i.e.,  $2V_D + V_{BAT}$ ).  $V_{BAT}$  can harness more energy when MOSFETs replace the diodes because the barrier is lower, but only after  $i_P$  charges  $C_P$  above  $V_{BAT}$ .



Fig. 9. Recycling inductor  $L_{RE}$  into a full-wave rectifier.

When unloaded, to be more specific,  $i_P$  charges  $C_P$  from negative to positive open-circuit voltages  $-V_{OC}$  to  $V_{OC}$  (by  $2V_{OC}$ ) with charge  $Q_{OC}$ , which is  $2V_{OC}C_P$ . When loaded, the rectifier conducts to  $V_{BAT}$  the portion of  $Q_{OC}$  that would have charged open-circuited  $C_P$  above  $|V_{BAT}|$  to  $|V_{OC}|$ , so  $V_{BAT}$  harnesses the difference twice (every half cycle) as

$$E_{\rm H} = 2(Q_{\rm OC} - 2V_{\rm BAT}C_{\rm P})V_{\rm BAT} = 4C_{\rm P}(V_{\rm OC} - V_{\rm BAT})V_{\rm BAT}, (6)$$

The peak of which happens at  $C_P V_{OC}^2$  when  $V_{BAT}$  is 0.5V<sub>OC</sub> [14]. Here, vibrations supply and absorb the energy with which they charge and discharge  $C_P$  between  $V_{BAT}$  and  $-V_{BAT}$ .

#### **Recycling Inductor:**

 $L_{RE}$  in Fig. 9 [15]–[16] increases  $E_H$  by recycling  $C_P$ 's energy at  $V_{BAT}$  to energize  $C_P$  in the other direction to  $-V_{BAT}$ . That is, after the positive half cycle,  $S_{RE}$  closes and  $L_{RE}$  de-energizes  $C_P$  and subsequently (through resonance) supplies the energy  $L_{RE}$  stored in the process to charge  $C_P$  to  $-V_{BAT}$ . In this manner,  $C_P$  draws no mechanical energy to charge to  $V_{BAT}$  and  $-V_{BAT}$ , so collects all of  $Q_{OC}$  as:

 $E_{\text{H}}' = 2Q_{\text{OC}}V_{\text{BAT}} = 2(2V_{\text{OC}}C_{\text{P}})V_{\text{BAT}} = 4C_{\text{P}}V_{\text{OC}}V_{\text{BAT}} \,. \label{eq:eq:energy_eq}$ 

 $S_{RE}$  and the circuit used to control  $S_{RE}$  dissipate power, so the energy  $C_P$  requires to charge between  $-V_{BAT}$  and  $V_{BAT}$  every half cycle, which is  $2(0.5C_P(2V_{BAT})^2)$  or  $4C_PV_{BAT}^2$ , should surpass these losses. In the end, drawing energy from vibrations amounts to damping them. With a rectifier, since the transducer ejects  $Q_{OC}$ near  $V_{BAT}$ , and output energy per half cycle is  $Q_{OC}V_{BAT}$ ,  $V_{BAT}$  ultimately limits the electrical damping force from which the transducer harvests energy.

#### **Reinvesting Energy [Proposed]:**

Increasing output energy is possible by reinvesting the energy gained in half the cycle (rather than depositing into  $V_{BAT}$ ) to increase the electrical damping force in the other half. For example, redirecting all the energy  $C_P$  draws from vibrations to charge by  $2V_{OC}$  to charge  $C_P$  in the opposite direction pre-charges  $C_P$  to  $-2V_{OC}$  so vibrations in the negative half cycle further charge  $C_P$  by another  $2V_{OC}$  to  $-4V_{OC}$ .



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Because the energy in a capacitor increases with the square of its voltage, harnessing what  $C_P$  stores at  $-4V_{OC}$  once per cycle produces more than drawing  $C_P$ 's energy twice at half that voltage, at  $2V_{OC}$  and  $-2V_{OC}$ :

 $E_{\rm H}^{\prime \prime \prime \prime} = 0 + 0.5 C_{\rm P} (-4 V_{\rm OC})^2 = 8 C_{\rm P} V_{\rm OC}^2$ .

To realize this, after  $CP_{PK}$  in Fig. 11 senses that  $v_P$  peaks,  $M_{N1}$ – $M_{N2}$  and  $M_{N3}$ – $M_{N4}$ , which implement  $S_I$  and  $S_N$  in Fig. 10, close for  $L_HC_P$ 's half resonance period so that  $C_P$  discharges into  $L_H$  and  $L_H$  subsequently de-energizes back into  $C_P$ . Once  $CP_{PK}$  senses that open-circuited  $C_P$  peaks in the opposite direction,  $S_I$  and  $S_N$  close to discharge  $C_P$  into  $L_H$  and  $S_I$  alone opens to de-energize  $L_H$  into  $V_{BAT}$  through  $M_{PDI}$ , which together with  $CP_{DI}$ , emulates diode  $D_I$ .



Fig. 11. Prototyped 2-µm BiCMOS re-investing, rectifier-free harvester.

#### **C. Experimental Validation**

The prototyped 2-µm BiCMOS harvester tunes the time that  $S_I$  and  $S_N$  connect ( $\tau_{REINV}$ ) externally with  $v_{REINV}$ ,  $\tau_{REINV}$  extends beyond  $L_HC_P$ 's quarter resonance period to a half to reinvest  $L_H$ 's energy back in  $C_P$ . When  $\tau_{REINV}$  is less than  $L_HC_P$ 's half resonance period,  $S_N$  opens early and  $L_H$  drains remnant energy into  $V_{BAT}$  via  $M_{PDN}$ , which with  $CP_{DN}$ , implements  $D_N$ . Once tuned, shaking a 44 × 13 × 0.4-mm<sup>3</sup> piezoelectric transducer charged  $C_P$  and  $L_H$  then recycled  $C_P$ 's energy at 1.02 V to pre-charge  $C_P$  in the opposite direction to -0.36 V, as Fig. 12a shows. Vibrations then charged  $C_P$  further to -1.9 V before  $L_H$  deenergized  $C_P$  into  $C_{BAT}$ . After 2.5 s of repeated cycles,  $C_{BAT}$  charged from 2.68 to 4.36 V, as Fig. 12b corroborates.

Without reinvesting energy,  $C_P$  charged to 1.4 V and – 1.2 V to energize  $C_{BAT}$  from 2.68 to 4.36 V in 3 s, which under similar conditions, means reinvesting energy produced 20% more output power.



Fig.12. Measured C<sub>P</sub> and C<sub>BAT</sub> charge profiles with and without reinvestment.

#### **D. Performance and Limitations**

Notice  $CP_{PK}$  is late in detecting  $v_P$ 's peaks, so before  $L_H$  can de-energize  $C_P$ , vibrations absorb some of  $C_P$ 's energy (in both cases shown). Also note that  $L_H$ 's reinvestment in  $C_P$  is unable to charge  $C_P$  to -1.02 V because conducting switches ( $S_N$  and  $S_I$ ) and  $L_H$ 's equivalent series resistance  $R_{ESR}$  dissipate some of that energy. This is critical because reducing  $C_P$ 's negative peak voltage has a squared impact on  $C_P$ 's peak energy, which is what the system harvests.



Fig. 13. Experimental and simulated output power across investment time.

Interestingly, as the experimental results of Fig. 13a show, increasing the investment in  $C_P$  produces diminishing returns in  $P_O$ . This results because transferring more energy through the switches and  $L_H$ 's  $R_{ESR}$  also increases conduction losses to the point they overwhelm reinvestment gains.



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Enlarging the FETs to lower their resistances balances losses and therefore raises  $P_0$ , as the simulated traces of Fig. 13b show. With twenty times (20×) larger FETs for  $S_N$  and  $S_I$  (at 72000µm/2µm), in fact, fully investing  $C_P$ 's positive energy into the negative phase raised simulated  $P_0$  by 56% from 47.4 µW to 74.2 µW. Ultimately, however, FET losses vary with input power, process, and temperature, but not mismatch.

### **V. CONCLUSIONS**

The experimental results of the piezoelectric transducer and the simulated results of the electromagnetic and electrostatic cases show that investing energy into the system increases output power Po. This is important because the coupling factors of tiny transducers and transponding inductors are substantially low, which means Po is also low. The idea here is to invest energy to raise the electrical damping force against which motion, magnetic fields, etc. work. This way, transducers draw more energy from the environment. The circuit components that transfer the investment, however, consume power, limiting the extent to which increased investments raise Po. Still, increasing Po this way, beyond reducing losses in the system, expands the functional reach of miniaturized systems to more practical levels.

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